Analysis of the Von Neumann stability criteria for a finite difference scheme for solving the 2D Riemannian wave equation

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Finite Difference Scheme for the Riemannanian 2D acoustic wave equation

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Finite Difference Scheme for the Riemannanian 2D acoustic wave equation

$$\begin{bmatrix} \nabla_{\xi}^{2} - \frac{1}{\nu_{\xi}^{2}} \frac{\partial^{2}}{\partial t^{2}} \end{bmatrix} U_{\xi} = F_{\xi}$$

$$\nabla_{\xi}^{2} = \frac{1}{\sqrt{|g|}} \frac{\partial}{\partial \xi_{i}} \left(g^{ij} \sqrt{|g|} \right) \frac{\partial}{\partial \xi_{j}} + g^{ij} \frac{\partial^{2}}{\partial \xi_{i} \partial \xi_{j}}$$

$$\nabla_{\xi}^{2} = \zeta^{i} \frac{\partial}{\partial \xi_{i}} + g^{ij} \frac{\partial^{2}}{\partial \xi_{i} \partial \xi_{j}}$$

► Finite Difference Scheme for the Riemannanian 2D acoustic wave equation Then, we have

$$\zeta^{i} \frac{\partial U_{\xi}}{\partial \xi_{i}} + g^{ij} \frac{\partial^{2} U_{\xi}}{\partial \xi_{i} \partial \xi_{j}} = \frac{1}{\nu_{\xi}^{2}} \frac{\partial^{2} U_{\xi}}{\partial t^{2}} + F_{\xi}$$

► Finite Difference Scheme for the Riemannanian 2D acoustic wave equation For a 2D scheme, we have

$$\frac{\partial^2 \textit{U}_\xi}{\partial t^2} = \nu^2 \left[\zeta^1 \frac{\partial \textit{U}_\xi}{\partial \xi_1} + \zeta^2 \frac{\partial \textit{U}_\xi}{\partial \xi_2} + \textit{g}^{11} \frac{\partial^2 \textit{U}_\xi}{\partial \xi_1^2} + 2 \textit{g}^{12} \frac{\partial^2 \textit{U}_\xi}{\partial \xi_1 \partial \xi_2} + \textit{g}^{22} \frac{\partial^2 \textit{U}_\xi}{\partial \xi^2} \right]$$

► Finite Difference Scheme for the Riemannanian 2D acoustic wave equation Take the following FD scheme

$$\begin{array}{lcl} \frac{\partial^{2} U}{\partial t^{2}} & = & \frac{U_{v,k}^{n+1} - 2U_{v,k}^{n} + U_{v,k}^{n-1}}{(\Delta t)^{2}} \\ \\ \frac{\partial U}{\partial \xi_{1}} & = & \frac{U_{v+1,k}^{n} - U_{v-1,k}^{n}}{2\Delta \xi_{1}} \\ \\ \frac{\partial U}{\partial \xi_{1}\partial \xi_{2}} & = & \frac{U_{v+1,k+1}^{n} - U_{v-1,k+1}^{n} - U_{v+1,k-1}^{n} + U_{v-1,k-1}^{n}}{2\Delta \xi_{1}\Delta \xi_{2}} \\ \\ \frac{\partial^{2} U}{\partial \xi_{1}^{2}} & = & \frac{U_{v+1,k}^{n} - 2U_{v,k}^{n} + U_{v-1,k}^{n}}{(\Delta \xi_{1}^{2})} \\ \\ \frac{\partial^{2} U}{\partial \xi_{2}^{2}} & = & \frac{U_{v,k+1}^{n} - 2U_{v,k}^{n} + U_{v,k-1}^{n}}{(\Delta \xi_{2}^{2})} \end{array}$$

where

$$\xi_1 = v\Delta\xi_1
\xi_2 = k\Delta\xi_2
t = n\Delta t
U_{v,k}^{\prime\prime} = U(\xi_1, \xi_2, t)$$

Finite Difference Scheme for the Riemannanian 2D acoustic wave equation
 So we obtain the followin discrete equation

$$\begin{split} U^n_{v,k} &=& 2U^n_{v,k} - U^{n-1}_{v,k} + (\nu \Delta t)^2 [\zeta^1 \left(\frac{U^n_{v+1,k} - U^n_{v-1,k}}{2\Delta \xi_1} \right) \\ &+ & \zeta^2 \left(\frac{U^n_{v,k+1} - U^n_{v,k-1}}{2\Delta \xi_2} \right) \\ &+ & g^{11} \left(\frac{U^n_{v+1,k} - 2U^n_{v,k} + U^n_{v-1,k}}{(\Delta \xi_1)^2} \right) \\ &+ & g^{22} \left(\frac{U^n_{v,k+1} - 2U^n_{v,k} + U^n_{v,k-1}}{(\Delta \xi_2)^2} \right) \\ &+ & g^{12} \left(\frac{U^n_{v+1,k+1} - U^n_{v-1,k+1} + U^n_{v+1,k-1} + U^n_{v-1,k-1}}{2\Delta \xi_1 \Delta \xi_2} \right)] \end{split}$$

Consider the exact solution to the Riemannian acoustic wave equation as

$$U(\xi_1,\xi_2,t)=U_{v,k}^n,$$

and the approximate solution as

$$u(\xi_1,\xi_2,t)=u_{v,k}^n.$$

The Von Neumann criteria states that

$$\varepsilon_{v,k}^n = U_{v,k}^n - u_{v,k}^n,$$

where $\varepsilon_{v,k}^n$ is the error of the wavefield at (ξ_1,ξ_2,t) . Suppose that we can decompose the error in terms of Fourier modes as

$$\varepsilon_{v,k}^{n} = e^{-i(\omega n \Delta t - k_{\xi_1} v \Delta \xi_1 - k_{\xi_2} k \Delta \xi_2)}$$

Inserting this aproximation into the discretized wave equation, we get

$$\begin{split} e^{-i\omega\Delta t}u^{n}_{v,k} &=& 2u^{n}_{v,k} - e^{-i\omega\Delta t}u^{n}_{v,k} \\ &+ & u^{n}_{v,k}(\nu\Delta t)^{2}\frac{\zeta^{1}}{2\Delta\xi_{1}}\left(e^{ik\xi_{1}\Delta\xi_{1}} - e^{-ik\xi_{1}\Delta\xi_{1}}\right) \\ &+ & \frac{\zeta^{2}}{2\Delta\xi_{2}}\left(e^{ik\xi_{2}\Delta\xi_{2}} - e^{-ik\xi_{2}\Delta\xi_{2}}\right) \\ &+ & \frac{g^{11}}{(\Delta\xi_{1})^{2}}\left(e^{ik\xi_{1}\Delta\xi_{1}} - e^{-ik\xi_{1}\Delta\xi_{1}} - 2\right) \\ &+ & \frac{g^{22}}{(\Delta\xi_{2})^{2}}\left(e^{ik\xi_{2}\Delta\xi_{2}} - e^{-ik\xi_{2}\Delta\xi_{2}} - 2\right) \\ &+ & \frac{g^{12}}{2\Delta\xi_{1}\Delta\xi_{2}}(e^{ik\xi_{1}\Delta\xi_{1}}e^{ik\xi_{2}\Delta\xi_{2}} \\ &- & e^{-ik\xi_{1}\Delta\xi_{1}}e^{ik\xi_{2}\Delta\xi_{2}} \\ &- & e^{ik\xi_{1}\Delta\xi_{1}}e^{-ik\xi_{2}\Delta\xi_{2}} + e^{-ik\xi_{1}\Delta\xi_{1}}e^{-ik\xi_{1}\Delta\xi_{1}} \end{split}$$

After some simplifications, we get

$$\Delta t \leq \frac{1}{\nu} \left(\frac{G-h}{(\Delta \xi)^2} - \frac{Z}{4\Delta \xi} \right)^{\frac{-1}{2}},$$

where

$$G = |g^{11} + g^{22} + 2g^{12}|$$

$$h = 2|g^{12}|$$

$$Z = |\zeta^{1} + \zeta^{2}|$$

Note that in the cartesian case $g_{ij} = \delta_{ij}$ we have

$$G = 2$$

$$h = 0$$

$$Z = 0$$

and then

$$\Delta t \leq \frac{\Delta \xi}{\nu \sqrt{2}}$$

which is the standar Courant stability criteria for FD 2D wave equation in Cartesian coordinates.

Some references

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