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# Wavelet as affine group, practical evidence and theoretical hypothesis

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# Outline

- Introduction
- Coherent States
- Reconstruction and decomposition
- Utility of the affine group
- Singularity Spectrum
- Practical Evidence
- Conclusions

# Introduction

# Coherent States

$$[U(a, b)f](x) = |a|^{-1/2} f\left(\frac{x-b}{a}\right)$$

$$U(a, b) U(a', b') = U(aa', b + ab')$$

$$\int_G d\mu(g) |\langle f, U(g)f \rangle|^2 < \infty$$

$$\int_G d\mu(g) \langle h_1, U(g)\tilde{f} \rangle \overline{\langle h_2, U(g)\tilde{f} \rangle} = C_{\tilde{f}} \langle h_1, h_2 \rangle$$

$$C_{\tilde{f}} = \langle A\tilde{f}, \tilde{f} \rangle$$

$$a^{-2} da db$$

$$(Af)^\wedge(\xi) = |\xi|^{-1} \hat{f}(\xi)$$

[1,2]



But this group structure was not exploited because she went to the discretely labeled wavelet families and these do not correspond to subgroups of the  $ax+b$  group.

Meanwhile, the mathematicians...

$$f = C_{\psi}^{-1} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{da db}{a^2} (T^{\text{wav}} f)(a, b) \psi^{a,b}$$

$$C_{\psi} = 2\pi \int_0^{\infty} d\xi |\xi|^{-1} |\hat{\psi}(\xi)|^2 = 2\pi \int_{-\infty}^0 d\xi |\xi|^{-1} |\hat{\psi}(\xi)|^2 < \infty$$

[14]

$$f = C_{\psi}^{-1} \int_0^{\infty} \frac{da}{a^2} \int_{-\infty}^{\infty} db T^{\text{wav}} f(a, b) \psi^{a,b}$$

But using a complex wavelet even for the analysis of real functions may have its advantages...

[14]

# Reconstruction and Decomposition

$$\psi_1, \psi_2 \in L^1(\mathbb{R})$$

$$\psi_2' \in L^2(\mathbb{R})$$

$$x\psi_2 \in L^1(\mathbb{R})$$

$$\hat{\psi}_1(0) = 0 = \hat{\psi}_2(0)$$

$$f = C_{\psi_1, \psi_2}^{-1} \int \frac{da}{a^2} \int db \langle f, \psi_1^{a,b} \rangle \psi_2^{a,b}$$

$$f(x) = C_{\psi_1, \psi_2}^{-1} \lim_{\substack{A_1 \rightarrow 0 \\ A_2 \rightarrow \infty}} \int_{A_1 \leq |a| \leq A_2} \frac{da}{a^2} \int_{-\infty}^{\infty} db \langle f, \psi_1^{a,b} \rangle \psi_2^{a,b}(x)$$

...

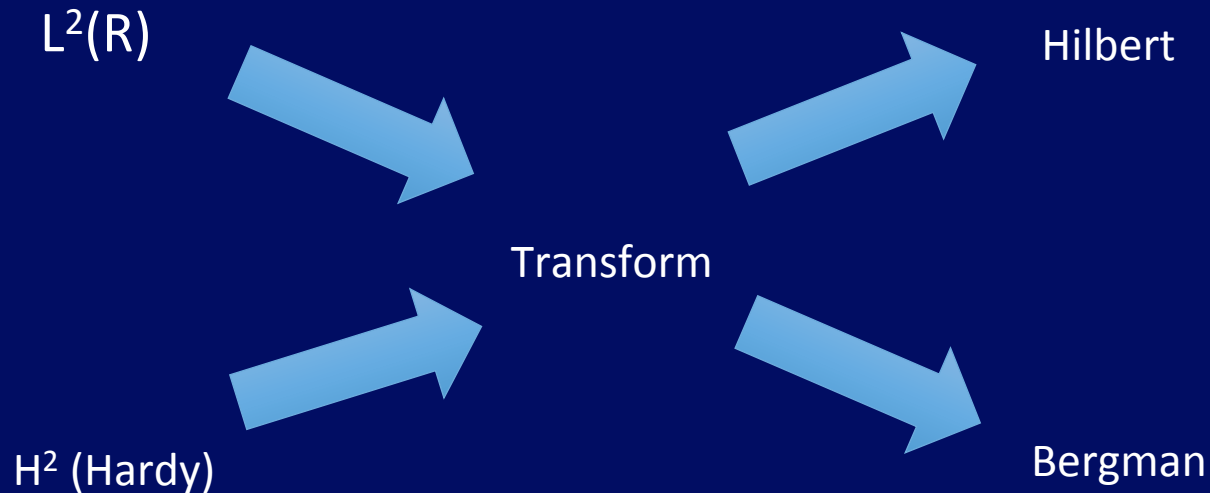
[14]

$$C_{\psi}^{-1} \int \int \frac{da db}{a^2} |(T^{\text{wav}} f)(a, b)|^2 = \int dx |f(x)|^2$$

The transform maps the space  $L^2(\mathbb{R})$  isometrically into a Hilbert space...

In some cases we can choose a subspace of  $L^2(\mathbb{R})$ , let's say  $H^2$  (Hardy)...  
The same transform can be interpreted as an isometry from  $H^2$  to the Bergman space

[14]



Then it is obvious that its values at certain discrete families of points completely determines the function... and this is sad

[14]

# Utility of the Affine Group

“... the utility of the affine group was possible because of the close connection between the affine and the canonical stories while the few distinctions can be used to advantage when appropriate...”

$$[Q, D] = i\hbar Q$$

- i) States of a system
- ii) Dynamics involves an automorphism among states
- iii) The natural inner product interpreted as probably amplitudes

OVERCOMPLETE FAMILY OF STATES

[3-8]



Indeed, modulo linear transformations, the affine variables form the elements of the only two-parameter, non-Abelian Lie algebra.

The affine group gets its name from its realization as the set of affine transformations of the real line given by

$$x \rightarrow x' = ax + b$$

Where

$$a \neq 0 \text{ and } b \in \mathbb{R}$$

# Singularity Spectrum

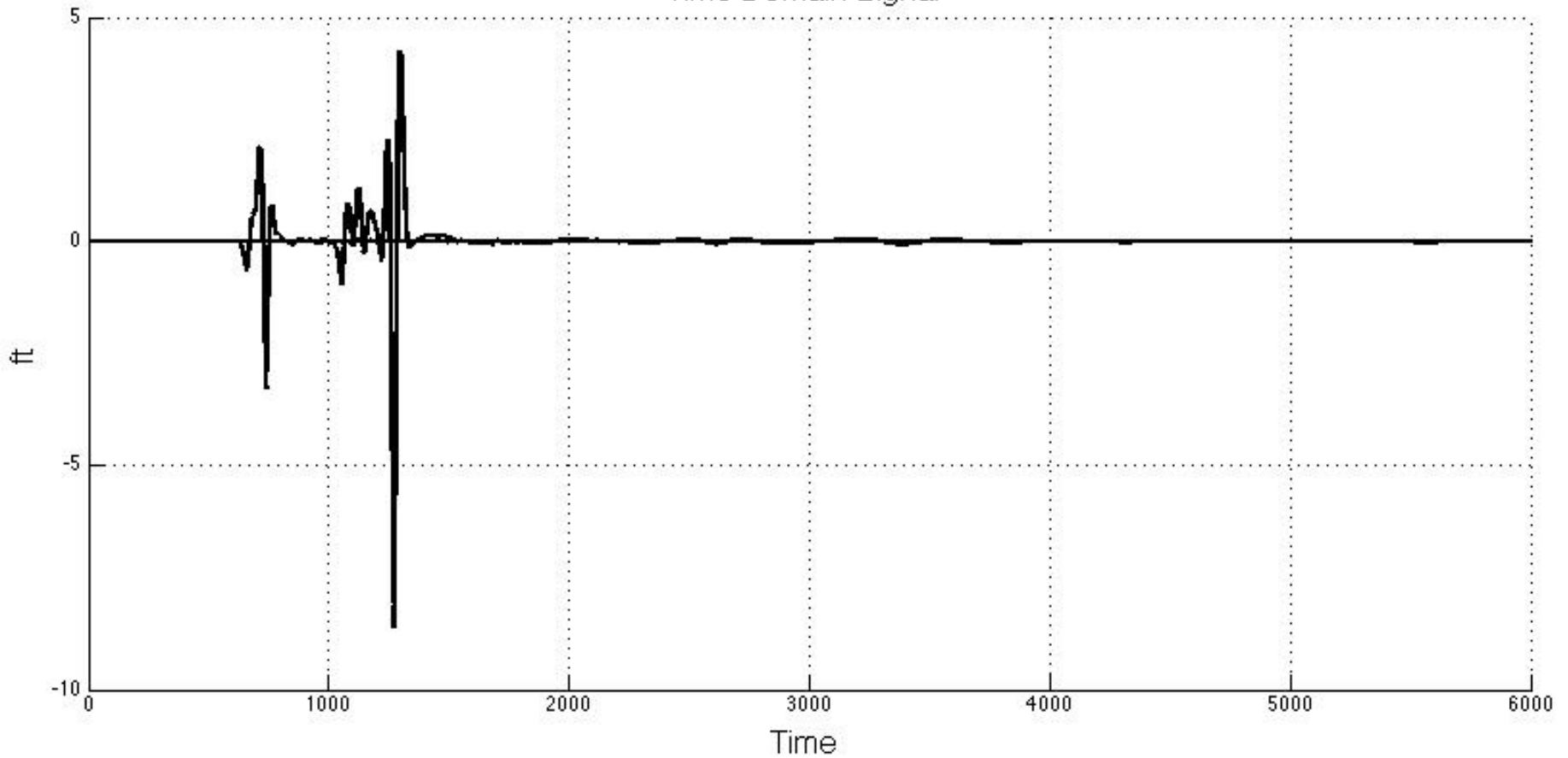
$$W_f(u, s) = \langle f(t), \psi_{u,s}(t) \rangle = \frac{1}{\sqrt{s}} \int_{-\infty}^{\infty} f(t) \psi^* \left( \frac{t-u}{s} \right) dt \quad f(t) = \frac{1}{C_\psi} \int_0^{\infty} \int_{-\infty}^{\infty} W_f(u, s) \frac{1}{\sqrt{s}} \psi \left( \frac{t-u}{s} \right) du \frac{ds}{s^2}$$

$$\int_{-\infty}^{\infty} t^p \psi(t) dt = 0, \quad p \in \{0, 1, \dots, (n-1)\}$$

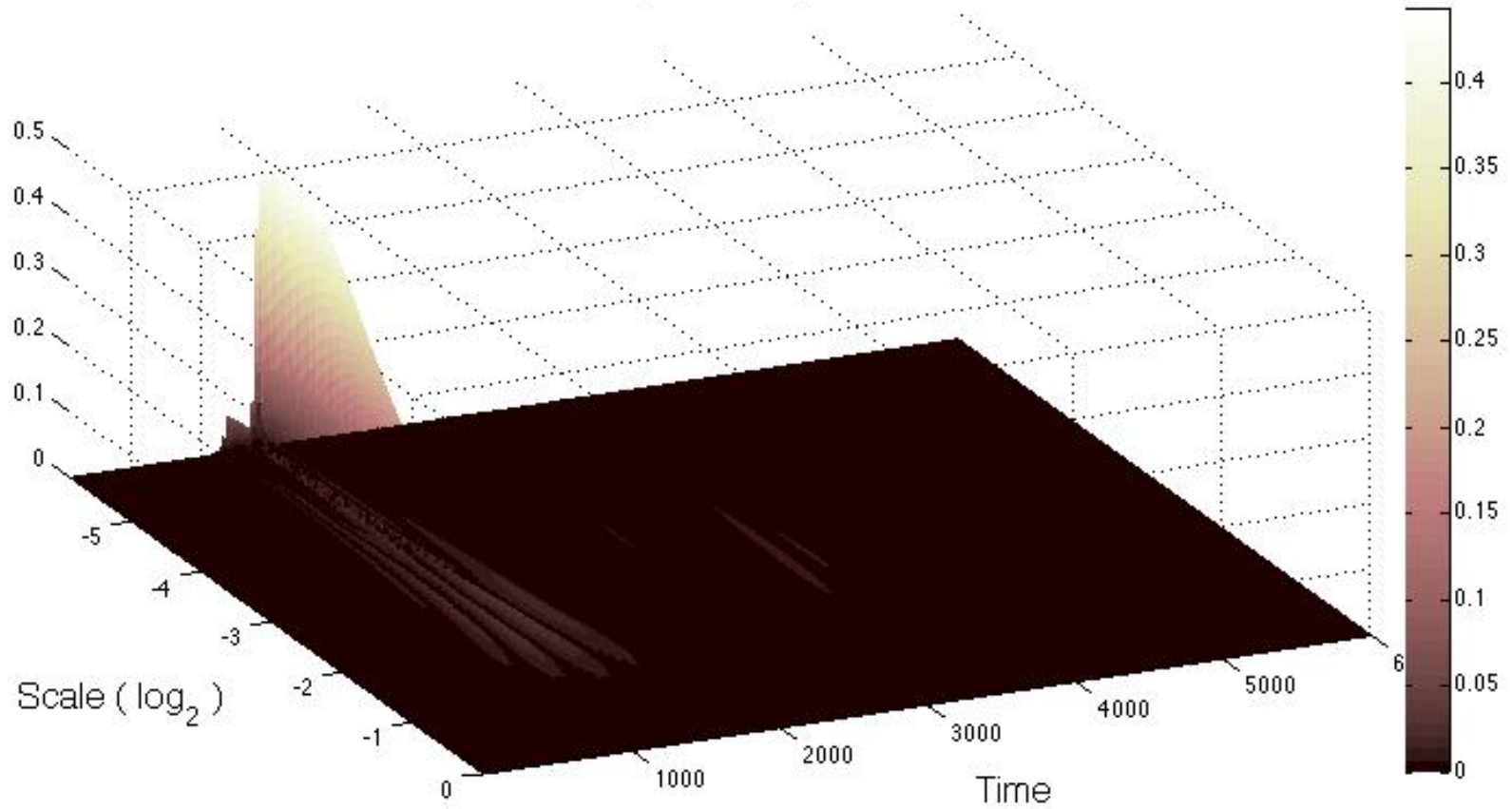
$$W_f(u, s) = f(u) * \bar{\psi}_s(u) \approx W_f[k, s_{i,j}] = \sum_{n=0}^{N-1} f[n] \cdot \bar{\psi}_{s_{i,j}}[n-k] = f[k] \star \bar{\psi}_{s_{i,j}}[k]$$

[12, 13, 14, 20, 21, 22]

Time Domain Signal



Scalogram of Signal ft



## Hölder Distribution

$$|f(t) - P_n(t - u)| \leq K|t - u|^\alpha$$

$$\alpha \geq 0 \quad C > 0, \quad \alpha \in (n, n + 1)$$

$$\int_{-\infty}^{\infty} |\hat{f}(\omega)|(1 + |\omega|^\alpha) d\omega < +\infty$$

Global regularity !

We want local !

$$|\psi(t)| \leq \frac{B_k}{1 + |t|^k} \quad \psi_n(t) = (-1)^n \frac{d^n \vartheta(t)}{dt^n}, \quad \vartheta(t) = e^{-t^2/2}$$

$$\bar{\psi}_{n,s}(t) = \frac{(-1)^n}{\sqrt{s}} \frac{d^n \vartheta(-t/s)}{dt^n} = \frac{1}{\sqrt{s}} s^n \frac{d^n \vartheta(\tau)}{d\tau^n} \Big|_{\tau=-t/s} = s^n \frac{d^n \bar{\vartheta}_s(t)}{dt^n}$$

[11, 15-17, 22-26]

## Hölder Exponent and Continuous Wavelet transform

$$W_{f,n}(u, s) = f(u) * \bar{\psi}_{n,s}(u) = f(u) * s^n \frac{d^n \bar{\vartheta}_s}{dt^n}(u) = s^n \frac{d^n}{dt^n} (f * \bar{\vartheta}_s)(u)$$

$$|W_f(u, s)| \leq Qs^{\alpha+1/2} \left( 1 + \left| \frac{u - \tau}{s} \right|^\alpha \right), \quad \forall (u, s) \in \mathbb{R} \times \mathbb{R}^+$$

$$|W_f(u, s)| \leq Qs^{\alpha+1/2} \left( 1 + \left| \frac{u - \tau}{s} \right|^\beta \right), \quad \forall (u, s) \in \mathbb{R} \times \mathbb{R}^+$$

[17, 18, 19]

## Wavelet transform Maximus Modulus

$$|W_f(u, s)| \leq Ms^{\alpha+1/2}$$

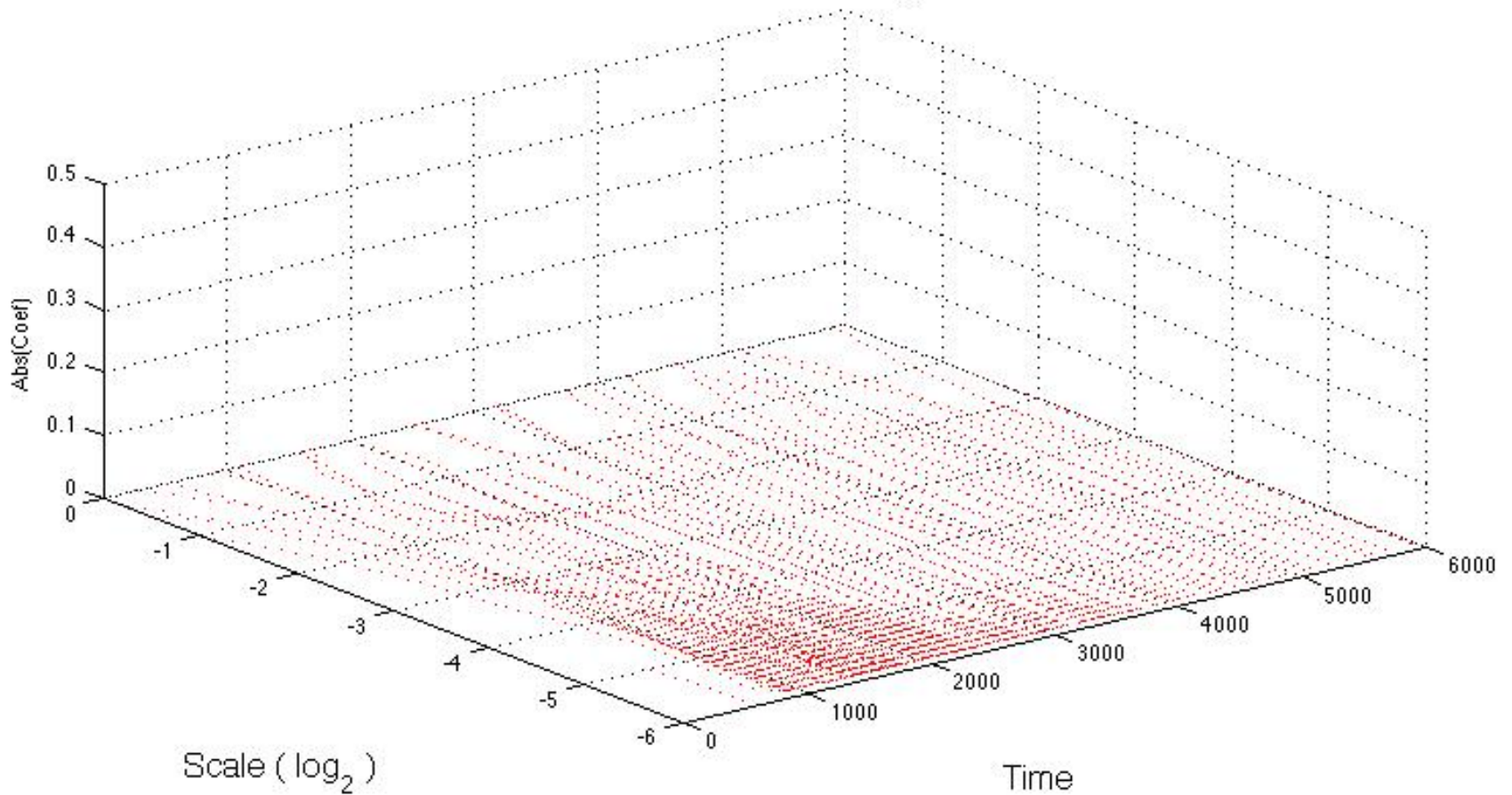
$$\log_2(|w_f(u, s)|) \approx \log_2(M) + \log_2(\alpha + 1/2) \log_2(s)$$

$$WTMM = \left\{ (u_0, s_0) \in (\mathbb{R}, \mathbb{R}^+), \frac{\partial |W_f(u, s)|}{\partial u} \Big|_{u=u_0, s=s_0} = 0 \right\}$$

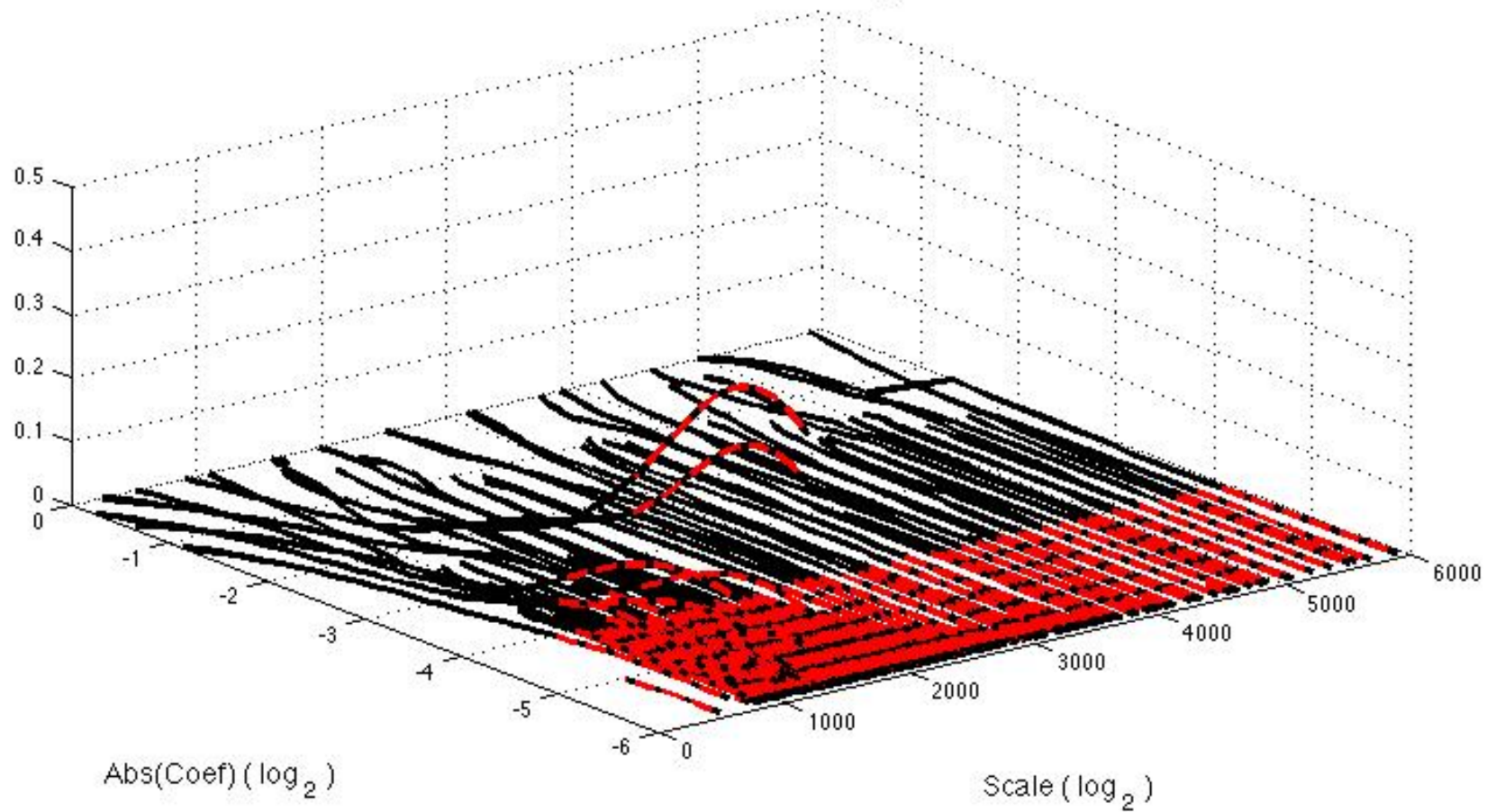
[20]

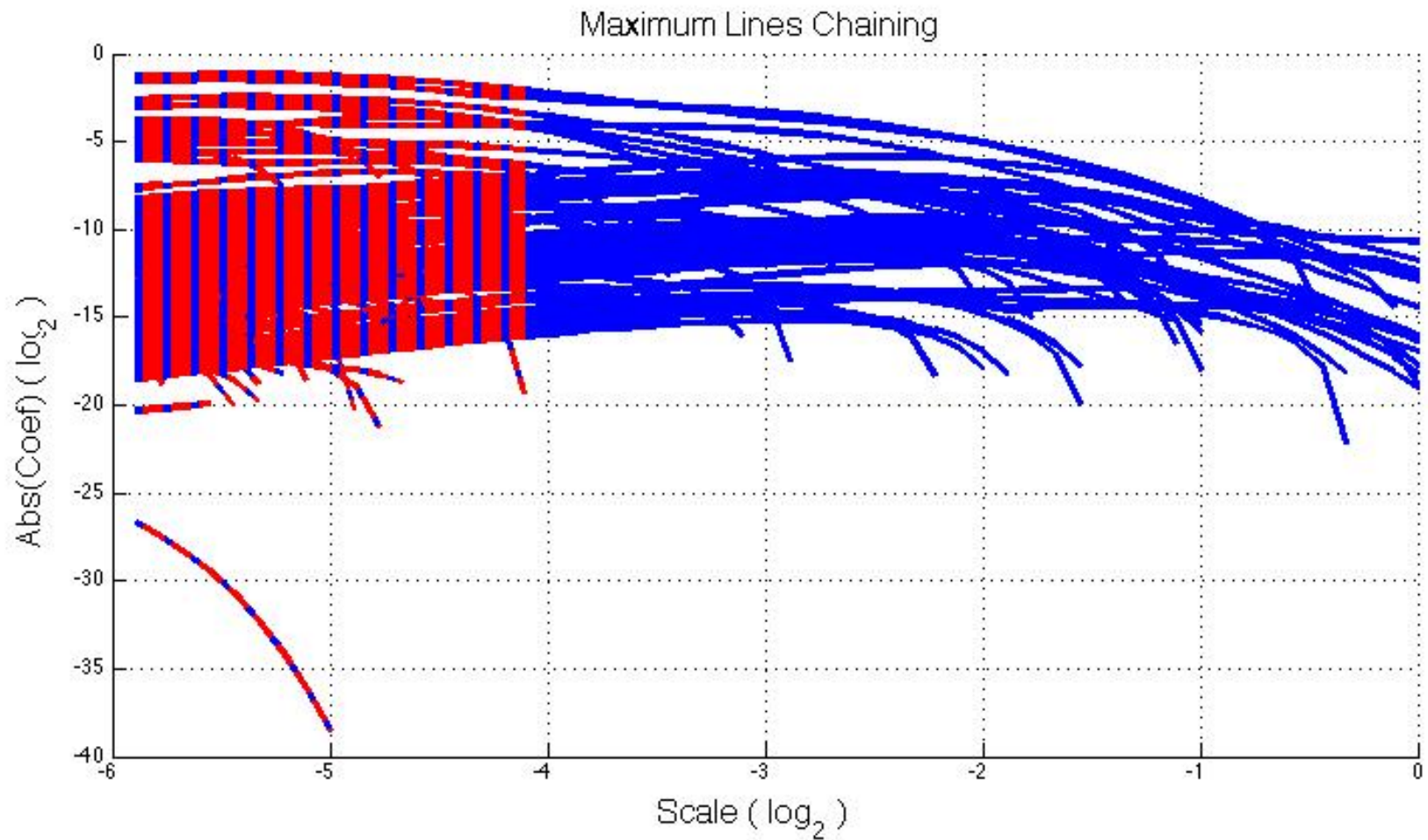


Maximum lines for CWT(ft)

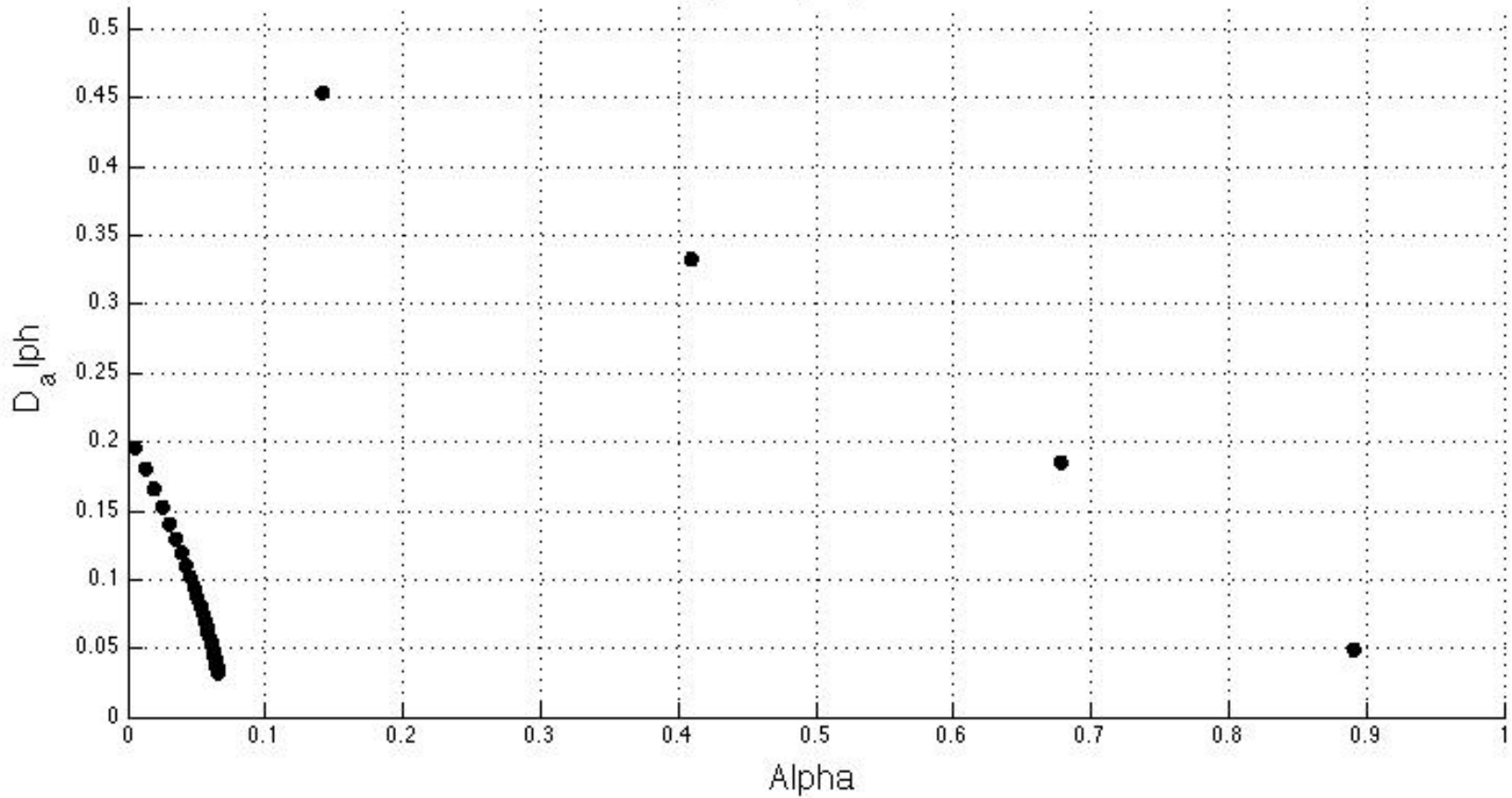


### Maximum Lines Chaining

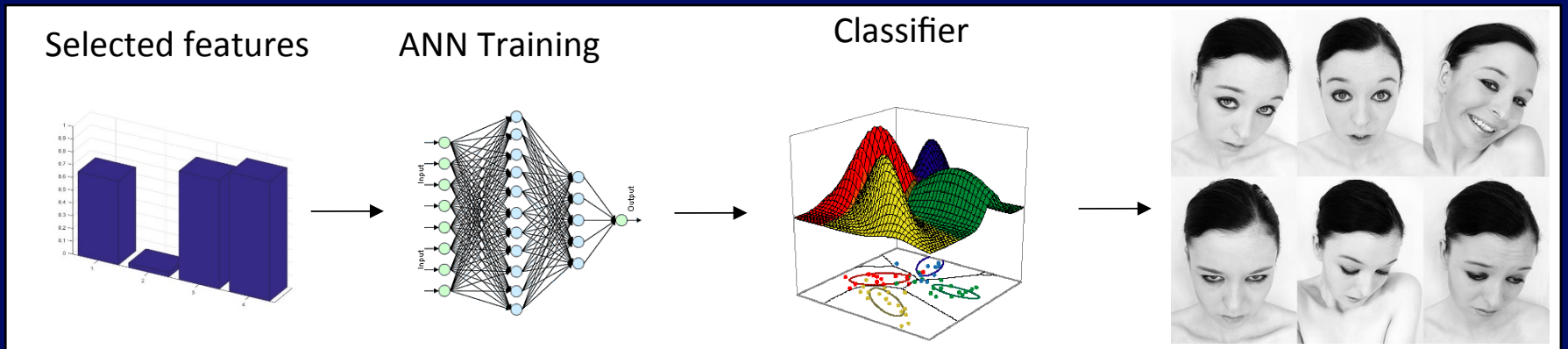
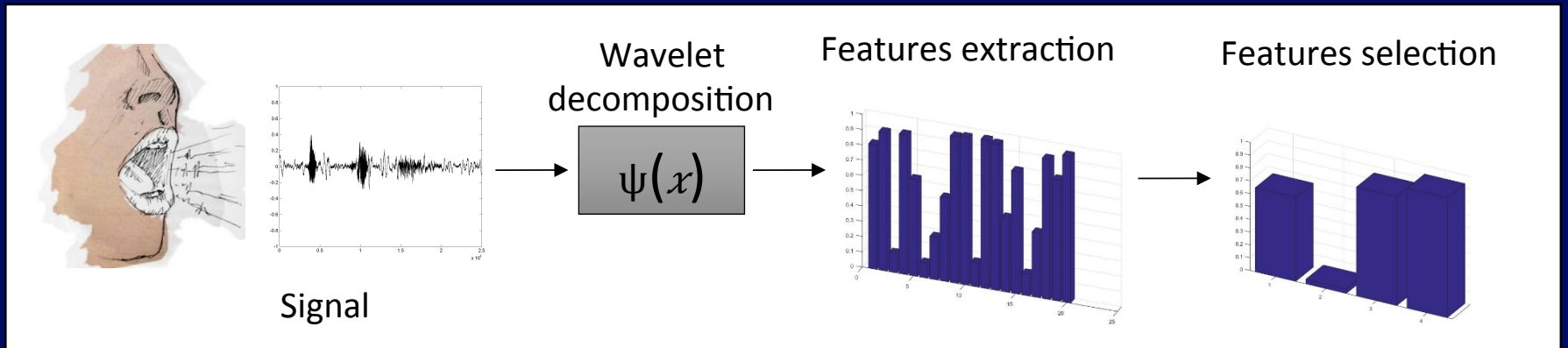




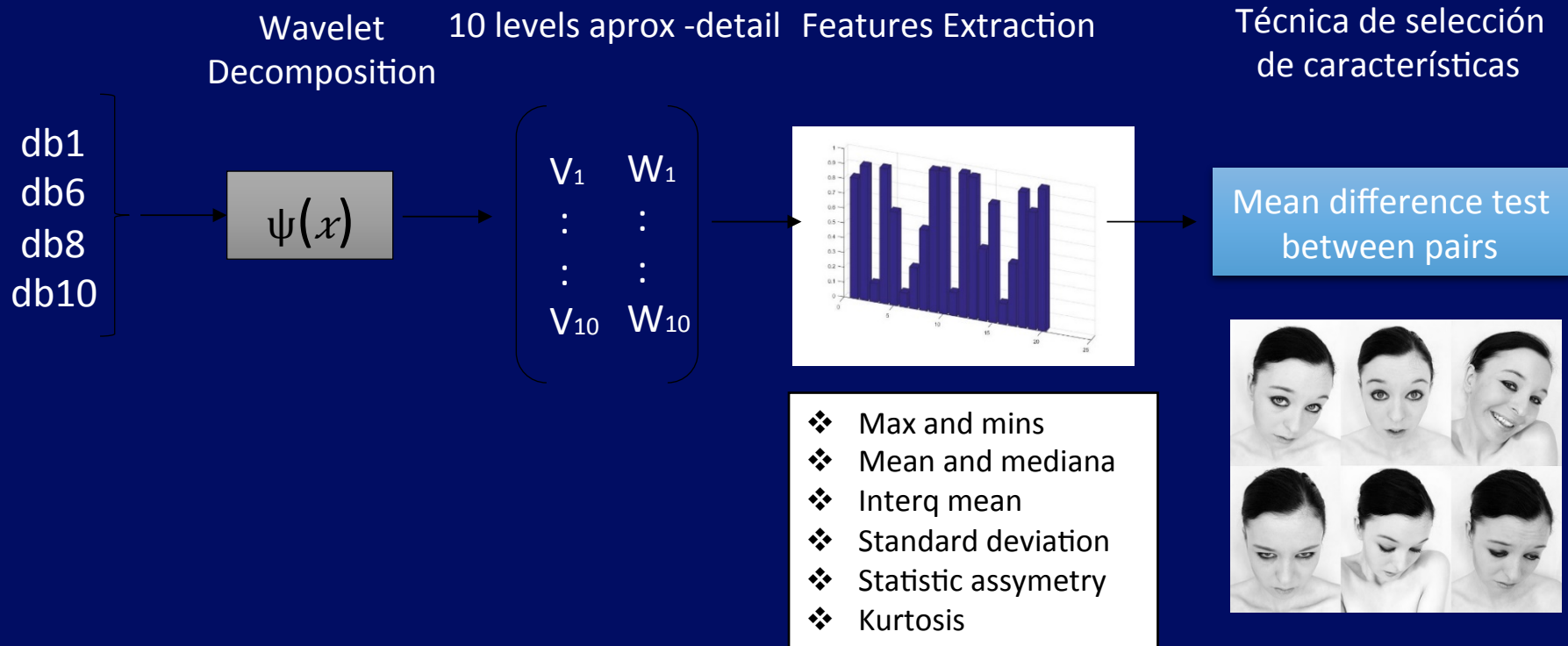
Singularity Spectrum



# Practical Evidence

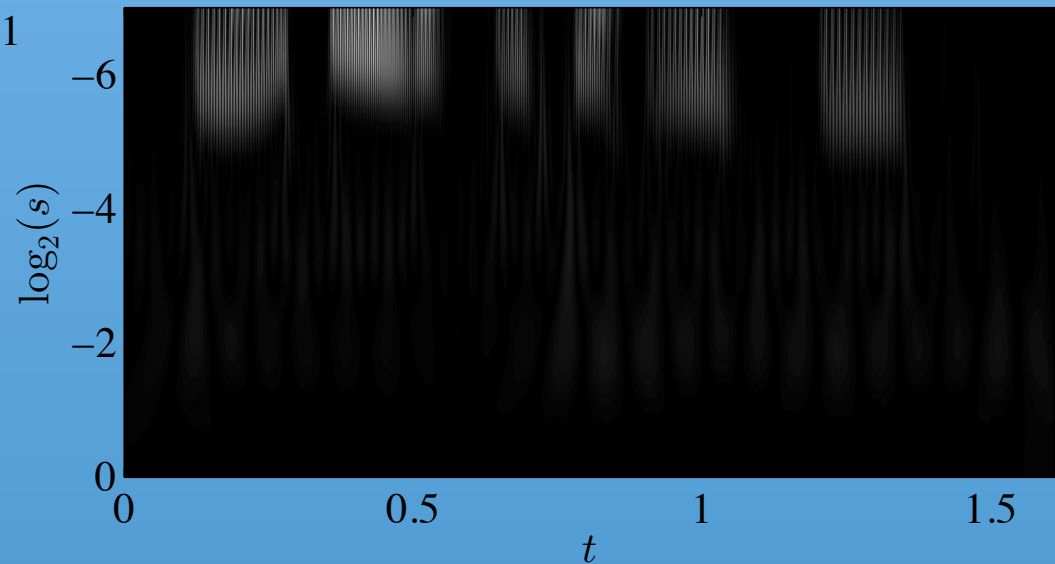
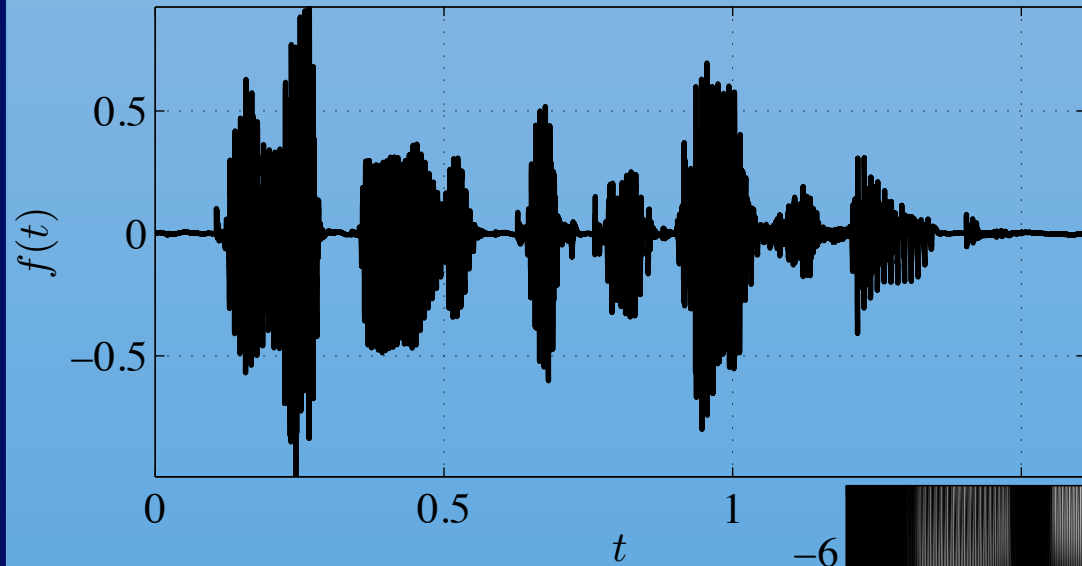


[27-30]



<b>EMOTIONS</b>	<i>BOREDOM</i>	<i>DISGUST</i>	<i>HAPPINESS</i>	<i>FEAR</i>	<i>NEUTRAL</i>	<i>ANGER</i>	<i>SADNESS</i>
<i>BOREDOM</i>	<b>89.74%</b>	0%	0%	1.28%	6.41%	1.28%	1.28%
<i>DISGUST</i>	2.22%	<b>91.11%</b>	4.44%	0%	2.22%	0%	0%
<i>HAPPINESS</i>	0%	2.86%	<b>85.71%</b>	2.86%	1.43%	7.14%	0%
<i>FEAR</i>	1.45%	1.45%	1.45%	<b>92.75%</b>	2.9%	0%	0%
<i>NEUTRAL</i>	6.25%	0%	0%	2.5%	<b>91.25%</b>	0%	0%
<i>ANGER</i>	0%	0.79%	4.76%	1.59%	0%	<b>92.86%</b>	0%
<i>SADNESS</i>	1.67%	0%	0%	1.67%	3.33%	0%	<b>93.33%</b>





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# Conclusions

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**— Thank you! —**